

Kepler 3:  $\omega_{1,K3}^2 r_1^3 = \omega_{2,K3}^2 r_2^3 = GM_0$ Spin-orbit resonance:  $\omega_{avg}^2 r_{avg}^3 = GM_0$ but also:  $\omega_{1,\text{actual}} = \omega_{\text{avg}} = \omega_{2,\text{actual}}$ as well as:  $r_1 < r_{avg} < r_2$ Therefore:  $\omega_{1,\text{actual}} < \omega_{1,\text{K3}}$ and:  $\omega_{2,actual} > \omega_{2,K3}$  $\Delta F_{cp} = \omega_{2,actual}^2 r_2 M_2 - \omega_{1,actual}^2 r_1 M_1 > 0$ We find: Roche limit:  $\Delta F_{q,1,2} = F_{q,0,1} - F_{q,0,2}$ Tear apart:  $\Delta F_{\text{tot}} = \Delta F_{a,1,2} + \Delta F_{\text{cp}} > \Delta F_{a,1,2}$ results in:  $M_2$  orbits too fast, goes off rails & gets in higher orbit  $M_1$  orbits too slow, dives inwards & gets in lower orbit as well as:  $M_1$  slows down more and more by  $M_0$ 's tidal bulge pulling it, and then: finally resulting in an impact (**IFF**  $M_0$  spins slower than  $\omega_1$ ). Innermost Nondestructive Circular Orbit  $\approx$  Roche limit. Conclusion: This is why every accretion disk has a hole in it. **ISCO? INCO!** 

In case of matter orbiting a black hole: replace *tidal bulges* with *frame dragging*. The latter is fully comparable to a *gravitational slingshot*,

see <a href="http://henk-reints.nl/astro/HR-frame-dragging.pdf">http://henk-reints.nl/astro/HR-frame-dragging.pdf</a> .

A slingshot of a spacecraft increases its tangential velocity. The centrifugal force (urh, shortage of centripetal force) then widens its (quasi) orbit around the Sun.

Similarly, the innermost matter of an accretion disk moves to a wider orbit, be it through the tidal effect or frame dragging.

Precondition is that the axial rotation of the central mass is faster than the other body's orbit.

p.3/3